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单自由度混联Ⅱ型惯容系统随机地震动响应分析^{*}

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摘要: 提出了基于 Clough-Penzien 谱的混联Ⅱ型惯容结构地震动响应封闭解的计算方法, 并基于所提方法研究了单自由度混联Ⅱ型惯容结构的减震性能及影响因素。首先, 建立混联Ⅱ型惯容结构的运动方程, 获得了结构位移与惯容出力等结构响应频域解的统一表达式。其次, 基于有理式分解与留数定理分别获得了频率响应特征值函数与 Clough-Penzien 谱的二次正交式, 进而获得了结构响应功率谱的二次正交式。最后, 得到了结构与惯容随机地震动响应 0~2 阶谱矩的简明封闭解。利用所提方法和虚拟激励法分别对一算例进行分析, 验证了所提方法的正确性; 同时, 利用所提方法分析了惯容各参数对结构抗震性能的影响。研究表明: 该文方法获得的响应解为封闭解, 在计算效率与计算准确性上都优于虚拟激励法。当 μ_ω 不变时, 随着 μ_m 与 μ_ξ 增大, 减震性能增强; 当 $\mu_\omega=1$ 时, 对减震性能的影响最大。

关 键 词: 惯容; Clough-Penzien 谱; 二次正交化; 谱矩; 封闭解

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Responses of SDOF Structures With SPIS-II Dampers Under Random Seismic Excitation

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Abstract: A closed-form solution of responses of SDOF structures with SPIS-II dampers under seismic excitation modeled with the Clough-Pezien spectrum was proposed, and the shock absorption performance and influential factors of this system were studied based on the proposed method. Firstly, the motion equation for the SPIS-II damper was established, and the unified expressions of frequency domain solutions of structural responses, such as the structural displacement and the inerter force, were obtained. Secondly, based on the rational expression decomposition and the residue theorem, the quadratic orthogonal equations of the frequency response eigenvalue function and the Clough-Pezien spectrum were obtained respectively, and in turn the quadratic orthogonal equation of the structural response power spectrum was deduced. Thirdly, the concise closed-form solutions of the 0~2nd-order spectral moments of the structural responses were acquired. The proposed method and the virtual excitation method were used to analyze a case respectively, which verifies the

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correctness of the proposed method. Finally, the proposed method was used to analyze the effects of the inerter parameters on the seismic performances of the structure. The research shows that, the proposed method gives closed-form solutions better than those given by the virtual excitation method in terms of computation efficiency and accuracy. The damping performance will improve with the increase of μ_m and μ_ξ for a constant μ_ω and the damping performance will reach the optimum for $\mu_\omega=1$.

Key words: inerter; Clough-Penzien spectrum; quadratic orthogonalization; spectral moment; closed-form solution

0 引言

被动控制技术^[1-2]是指在主结构的适当部位增加耗能装置或子结构系统而成的耗能体系,在外部激励作用下主结构带动耗能装置(子结构)运动,从而达到减小主结构振动的目的,其具有性能稳定、安装方便等优点。调谐质量阻尼器(TMD)^[3]、调液质量阻尼器(TLD)^[4]、隔震垫^[5]、黏弹性阻尼器^[6]等均为被动控制技术中成熟应用的装置,实践表明,它们在土木工程防震减灾中发挥了重要作用。新型被动控制装置的研发一直是工程师们追逐的目标,近年来,惯容系统^[7-10]作为一种新型被动控制装置,具有良好的减震性能,得到了广泛的关注。

Nakamura 和 Kawamata 等^[11]最早提出运用液体泵来减小结构的地震动响应,之后又提出利用加速液态质量阻尼器^[11]来减少建筑结构的地震动,两种装置均具有良好的减震性能,它们均是利用惯容原理制作的阻尼器,是惯容减震的最初形式。Arakaki 等^[12]提出将滚珠丝杠与黏滞阻尼器结合,使黏滞阻尼器所产生的阻尼力放大数倍,从而提高了阻尼器的减震效果,该方法是利用滚珠丝螺杆机制获得惯容法的应用。Saito 等^[13-14]设计出具有惯容特性的减震装置(滚珠丝杠振动控制装置),并对惯容增效与阻尼增效现象开始了研究。Saito 等^[15]提出调谐黏滞质量阻尼器(TVMD)并和 TMD 系统进行了比较,证明了对于带有阻尼器且附加阻尼系数相同的线性结构系统,TVMD 的振动控制是有效的。Pan 和 Zhang^[16]不仅证明了惯容系统的耗能增效方程,而且揭示了耗能增效的本质,同时将惯容系统应用于建筑结构的振动控制中。研究表明:设置惯容系统可有效降低结构位移、基层剪力和弯矩,建立了惯容系统与振动控制效果之间的桥梁。吴应雄等^[17]对长周期地震动作用下惯容-层间隔震结构响应进行时程分析,提出了可便捷有效地确定惯容系统参数的设计方法,并证明了混联Ⅱ式惯容结构对减小隔震层位移具有良好作用。上述研究表明,惯容系统拥有良好的减震性能。

大量研究表明,地震动具有显著的随机性,为此,工程师们先后提出白噪声^[18]、Kanai-Tajimi 谱^[19-20]和 Clough-Penzien 谱^[21]等随机地震动模型,模型参数越多,越能准确描述实际地震动。Zhang 等^[22]研究了惯容单自由度系统基于白噪声激励下的响应方差,给出了其封闭解。赵志鹏等^[23]运用能量法对基于白噪声激励的惯容系统减震性能与耗能效果进行分析,并通过算例验证了惯容系统能有效耗散结构的振动能量。潘超等^[24]分析了单自由度混联Ⅱ型惯容减震体系基于 Kanai-Tajimi 谱在地震作用下的实用设计方法,给出了结构响应方差的求解方法,但该方法需要对频率响应函数模值的平方在圆频率 $[0, +\infty]$ 积分。Clough-Penzien 谱是在 Kanai-Tajimi 谱的基础上增加一次滤波,解决了零频处奇异点的缺陷,更符合实际地震的随机激励模型。Ge 等^[25-26]运用功率谱的二次正交化法研究了 Kanai-Tajimi 谱与 Davenport 风速谱下结构系列响应简明封闭解,研究表明二次正交化法是一种研究平稳激励下结构动力响应谱矩和方法的新解法。

本文针对混联Ⅱ型惯容单自由度系统,首先,推导出混联Ⅱ型惯容微分型本构关系表达式,并重构了该系统的运动方程;其次,通过复模态对重构后的运动方程进行解耦,并对结构动态响应的频域解进行了统一化;然后,基于文献[25-26]的方法,对频率响应特征值函数进行二次正交化,再利用留数定理,对 Clough-Penzien 谱功率谱密度函数进行二次正交化,并与二次正交化后的频率响应特征值函数相结合,得到了混联Ⅱ型结构动态响应的功率谱密度函数正交化后的简明表达式,进而获得结构响应谱矩的简明封闭解;最后,通过具体算例来验证计算结果并对主要参数进行分析。

1 运动方程的重构

单自由度体系中设置惯容减震系统^[7],其计算简图如图 1 所示,地震激励下的运动方程为

$$m\ddot{x} + c\dot{x} + kx + F_{IS} = -m\ddot{x}_g, \quad (1)$$

式中, m , c , k 分别为单自由度结构质量、阻尼、刚度; x , \dot{x} , \ddot{x} 分别为主结构位移、速度、加速度; \ddot{x}_g 为地面运动加速度; F_{IS} 为惯容系统出力。

混联 II 型惯容系统^[7] 是由惯容与阻尼器并联后再与弹簧串联而成, 计算简图如图 2 所示, 有如下力学关系:

$$F_{IS} = k_s x_1 = c_d \dot{x}_2 + m_{in} \ddot{x}_2, \quad (2a)$$

$$x = x_1 + x_2, \quad (2b)$$

式中, k_s 为惯容系统中弹簧的刚度; m_{in} 为惯容质量; c_d 为惯容阻尼; x_2 为惯容系统阻尼元件与惯容元件并联部分的相对位移; x_1 为惯容系统中弹簧两端相对位移。

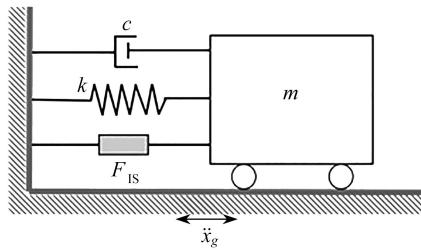


图 1 结构计算简图

Fig. 1 The structural calculation diagram

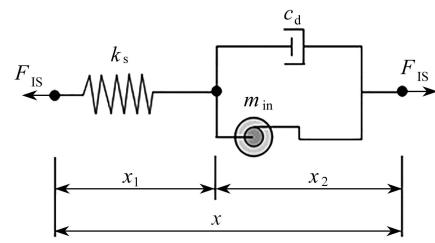


图 2 惯容系统计算简图

Fig. 2 The calculation diagram of the inerter system

由式(2)可知, 混联 II 型惯容系统中存在如下力学关系:

$$\dot{x} = \dot{x}_1 + \dot{x}_2, \quad (3a)$$

$$x_1 = F_{IS}/k_s. \quad (3b)$$

将式(3)代入式(2), 式(2)代入式(1)分别得

$$k_s x_2 + c_d \dot{x}_2 + m_{in} \ddot{x}_2 = k_s x, \quad (4a)$$

$$m\ddot{x} + c\dot{x} + kx + c_d \dot{x}_2 + m_{in} \ddot{x}_2 = -m\ddot{x}_g. \quad (4b)$$

联立式(4a)和式(4b), 并用状态方程表示:

$$M\dot{X} + KX = \alpha\ddot{x}_g, \quad (5)$$

$$X = [x \ x_2 \ \dot{x} \ \dot{x}_2]^T, \quad (6)$$

式中, $\alpha = [-m \ 0 \ 0 \ 0]^T$, 上标“T”表示矩阵转置,

$$M = \begin{bmatrix} c & 0 & m & m_{in} \\ 0 & c_d & 0 & m_{in} \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix}, \quad K = \begin{bmatrix} k & 0 & 0 & c_d \\ -k_s & k_s & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}.$$

2 结构地震动响应频域解的统一表达式

根据复模态理论^[27-29], 方程(5)存在对角阵 P 和左、右特征向量 V, U 使之解耦, 即存在如下关系:

$$P = \frac{V^T K U}{V^T M U}. \quad (7)$$

引入复模态变换

$$X = Uz, \quad (8)$$

式中, z 为复模态广义参数。将式(8)代入式(5), 并左乘 V^T , 可得

$$\dot{z} + \frac{V^T K U}{V^T M U} z = \eta \ddot{x}_g, \quad (9)$$

式中 $\eta = V^T \alpha / (V^T M U)$.

将式(7)代入式(9), 则式(9)可以化简为

$$\dot{z} + p z = \eta \ddot{x}_g. \quad (10)$$

由于 p 为对角阵, 则式(10)的分量形式为

$$\dot{z}_k + p_k z_k = \eta_k \ddot{x}_g, \quad (11)$$

其中, z_k, η_k, p_k 分别为 z, η, p 的分量.

由式(7)和式(11)可知 η_k 与 p_k 只与惯容系统的力学参数相关, 而与地震激励无关. 当惯容系统为线性时, 上述参数易于求解. 为便于后文应用, p_k 与 η_k 称之为系统模态特征值和系统模态强度系数.

虚拟激励法^[30-32]由林家浩于 20 世纪 80 年代提出, 其特点是将随机激励转化为简谐激励, 进而获得结构响应功率谱密度函数的精确表达式. 针对式(11)引入虚拟激励荷载

$$\ddot{x}_g(\omega, t) = \sqrt{S_{\ddot{x}_g}(\omega)} e^{i\omega t}, \quad (12)$$

式中, $i = \sqrt{-1}$, $S_{\ddot{x}_g}(\omega)$ 为地震加速度 $\ddot{x}_g(t)$ 的功率谱密度函数.

将式(12)代入式(11), 由虚拟激励法, 则

$$\bar{z}_k(\omega) = \frac{\eta_k}{p_k + i\omega} \sqrt{S_{\ddot{x}_g}(\omega)} e^{i\omega t}, \quad (13)$$

式中, $\bar{z}_k(\omega)$ 为式(11)的频域解.

由式(6)及式(11), 可得惯容减震系统的位移与速度频域解为

$$x(\omega) = \sum_{k=1}^4 u_{1,k} \bar{z}_k(\omega), \quad (14a)$$

$$\dot{x}(\omega) = \sum_{k=1}^4 u_{3,k} \bar{z}_k(\omega), \quad (14b)$$

$$x_2(\omega) = \sum_{k=1}^4 u_{2,k} \bar{z}_k(\omega), \quad (14c)$$

$$\dot{x}_2(\omega) = \sum_{k=1}^4 u_{4,k} \bar{z}_k(\omega). \quad (14d)$$

式中, $u_{l,k}$ 为右特征向量 \mathbf{U} 第 l 行第 k 列的元素.

由式(2)、(14c)及式(14d), 可得惯容出力的频域解为

$$F_{IS} = \sum_{k=1}^4 k_s (u_{1,k} - u_{2,k}) \bar{z}_k(\omega). \quad (15)$$

至此, 结构位移、速度与惯容出力的频域解可统一表示为

$$X_l(\omega) = \sum_{k=1}^4 \kappa_{l,k} \bar{z}_k(\omega), \quad (16)$$

式中, $X_l(\omega)$ 表示响应量 l 的频域解; $\kappa_{l,k}$ 表示响应量 l 的第 k 个模态值对应的系数, 称之为响应模态系数. 由式(14)–(16)可知, $\kappa_{l,k}$ 与系统的右特征向量 \mathbf{U} 及惯容参数有关. 而右特征向量 \mathbf{U} 又与整个惯容系统的力学参数有关, 其由复模态法易于求出.

3 响应功率谱的二次正交式

3.1 频率响应特征值函数的二次正交式

由虚拟激励法^[30-32]及式(15), 结构响应 X_l 的功率谱为

$$S_{X_l}(\omega) = \sum_{k=1}^2 \sum_{j=1}^2 \kappa_{l,k} \kappa_{l,j} \bar{z}_k(\omega) \bar{z}_j^*(\omega), \quad (17)$$

式中, $\bar{z}_j^*(\omega)$ 为 $\bar{z}_j(\omega)$ 的复共轭项.

将式(13)代入式(17)可得

$$S_{X_l}(\omega) = \sum_{k=1}^4 \sum_{j=1}^4 \kappa_{l,k} \frac{\eta_k}{p_k + i\omega} \sqrt{S_{\ddot{x}_g}(\omega)} e^{i\omega t} \kappa_{l,j} \frac{\eta_j}{p_j - i\omega} \sqrt{S_{\ddot{x}_g}(\omega)} e^{-i\omega t}. \quad (18)$$

式(18)可简化为

$$S_X(\omega) = S_{\ddot{x}_g}(\omega) H_X(\omega), \quad (19)$$

式中

$$H_X(\omega) = \sum_{k=1}^4 \sum_{j=1}^4 \kappa_k \eta_k \kappa_j \eta_j \frac{1}{p_k + i\omega} \frac{1}{p_j - i\omega}. \quad (20)$$

根据式(19)可以看出, 响应功率谱为 Clough-Penzien 谱地震动激励功率密度函数 $S_{\ddot{x}_g}(\omega)$ 与 $H_X(\omega)$ 的乘积. 由随机振动理论^[27]中的频域法可知, 结构响应功率谱密度函数等于激励的功率密度函数与结构频响函数模值的平方的乘积, 故 $H_X(\omega)$ 在数值上为结构频率响应函数的模值的平方, 称之为频率响应特征值函数^[28]. 根据文献[26]中的方法, 对 $H_X(\omega)$ 做如下等价变化:

$$H_X(\omega) = \sum_{j=1}^4 (\kappa_j \eta_j)^2 \frac{1}{p_j + i\omega} \frac{1}{p_j - i\omega} + \sum_{k=1}^3 \sum_{j=k+1}^4 \kappa_k \eta_k \kappa_j \eta_j \left(\frac{1}{p_k + i\omega} \frac{1}{p_j - i\omega} + \frac{1}{p_j + i\omega} \frac{1}{p_k - i\omega} \right), \quad (21)$$

式(20)进一步简化为

$$H_X(\omega) = \sum_{j=1}^4 \frac{(\kappa_j \eta_j)^2}{p_j^2 + \omega^2} + 2 \sum_{k=1}^3 \sum_{j=k+1}^4 \frac{\kappa_k \eta_k \kappa_j \eta_j}{p_k + p_j} \left(\frac{p_j}{p_j^2 + \omega^2} + \frac{p_k}{p_k^2 + \omega^2} \right). \quad (22)$$

3.2 Clough-Penzien 谱的二次正交式

Clough-Penzien 谱激励的功率谱密度函数为

$$S_{\ddot{x}_g}(\omega) = \frac{\omega^4}{(\omega_h^2 - \omega^2)^2 + 4\xi_h^2 \omega_h^2 \omega^2} \frac{\omega_g^4 + 4\xi_g^2 \omega_g^2 \omega^2}{(\omega_g^2 - \omega^2)^2 + 4\xi_g^2 \omega_g^2 \omega^2} S_0, \quad (23)$$

式中, ξ_g 和 ω_g 分别为场地土(第一滤波)的阻尼比和卓越频率; ξ_h 和 ω_h 分别为场地土(第二滤波)的阻尼比和卓越频率.

针对地震动功率谱密度函数, 利用留数定律^[33], 可获得其等效的二次正交式:

$$S_{\ddot{x}_g}(\omega) = S_0 \sum_{r=1}^2 \sum_{k=1}^2 \left(1 - \frac{c_{h,k}}{\omega^2 + \omega_{h,k}^2} \right) \frac{c_{g,r}}{\omega^2 + \omega_{g,r}^2}, \quad (24)$$

式中

$$\bar{\omega}_{h,1}^2 = (1 - 2\xi_h^2) \omega_h^2 + 2\omega_h^2 \xi_h \sqrt{\xi_h^2 - 1}, \quad \bar{\omega}_{h,2}^2 = \text{conj}(\bar{\omega}_{h,1}^2), \quad \omega_{h,k}^2 = -\bar{\omega}_{h,k}^2 \quad (k = 1, 2),$$

$$\bar{\omega}_{g,1}^2 = (1 - 2\xi_g^2) \omega_g^2 + 2\omega_g^2 \xi_g \sqrt{\xi_g^2 - 1}, \quad \bar{\omega}_{g,2}^2 = \text{conj}(\bar{\omega}_{g,1}^2), \quad \omega_{g,r}^2 = -\bar{\omega}_{g,r}^2 \quad (r = 1, 2),$$

$$c_{h,1} = \frac{\omega_h^4 + 2\omega_h^2 \omega_{h,1}^2 (2\xi_h^2 - 1)}{\bar{\omega}_{h,1}^2 - \bar{\omega}_{h,2}^2}, \quad c_{g,1} = \frac{\omega_g^4 + 4\xi_g^2 \omega_g^2 \omega_{g,1}^2}{\bar{\omega}_{g,1}^2 - \bar{\omega}_{g,2}^2},$$

$$c_{h,2} = \text{conj}(c_{h,1}), \quad c_{g,2} = \text{conj}(c_{g,1}),$$

其中“conj”表示复共轭项.

3.3 结构地震动响应功率谱密度函数的二次正交式

由式(19)、(22)及(24), 则结构地震动响应的功率谱可表示为

$$S_X(\omega) = S_0 \sum_{j=1}^4 (\kappa_j \eta_j)^2 A_j(\omega) + 2S_0 \sum_{k=1}^3 \sum_{j=k+1}^4 \frac{\kappa_k \eta_k \kappa_j \eta_j}{p_j + p_k} (p_j A_j(\omega) + p_k A_k(\omega)), \quad (25)$$

式中

$$A_j(\omega) = S_0 \sum_{r=1}^2 \sum_{k=1}^2 \frac{c_{g,r}}{\omega^2 + \omega_{g,r}^2} \left(1 - \frac{c_{h,k}}{\omega^2 + \omega_{h,k}^2} \right) \frac{1}{p_j^2 + \omega^2}. \quad (26)$$

由式(25)及(26)可知,结构地震动响应功率谱密度函数转化为地震动卓越频率与结构频率响应特征值函数的简明公式,为后文结构响应方差和0~2阶谱矩的分析奠定了基础。

4 地震动响应0~2阶谱矩的封闭表达式

由谱矩定义^[24],地震动响应的 q 阶谱矩为

$$\alpha_{X,q} = 2 \int_0^\infty \omega^q S_X(\omega) d\omega, \quad q = 0, 1, 2. \quad (27)$$

4.1 地震动响应的0阶谱矩及1阶谱矩简明封闭解

将式(26)代入式(27),则结构响应的谱矩为

$$\alpha_{X,q} = 2S_0 \left[\sum_{j=1}^4 (\kappa_j \eta_j)^2 T_j^q + 2 \sum_{j=1}^3 \sum_{k=j+1}^4 \frac{\kappa_k \eta_k \kappa_j \eta_j}{p_k + p_j} (p_j T_j^q + p_k T_k^q) \right], \quad (28)$$

式中

$$T_j^q = \int_0^\infty \omega^q A_j(\omega) d\omega. \quad (29)$$

将 $q=0$ 及 $q=1$ 分别代入式(29)中,可获得 T_j^0, T_j^1 的值,具体推导过程见附录。

$$T_j^0 = \frac{\pi}{2} \left\{ \sum_{r=1}^2 c_{g,r} [(p_j + \omega_{g,r}) \omega_{g,r} p_j]^{-1} - \sum_{r=1}^2 \sum_{k=1}^2 \frac{c_{h,k} c_{g,r}}{\omega_{h,k}^2 - \omega_{g,r}^2} \left[(\omega_{g,r} p_j^2 + \omega_{g,r}^2 p_j)^{-1} - (\omega_{h,k} p_j^2 + \omega_{h,k}^2 p_j)^{-1} \right] \right\} S_0, \quad (30a)$$

$$T_j^1 = \frac{1}{2} \left\{ \sum_{r=1}^2 c_{g,r} (\ln p_j^2 - \ln \omega_{g,r}^2) (p_j^2 - \omega_{g,r}^2)^{-1} - \sum_{r=1}^2 \sum_{k=1}^2 \frac{c_{h,r} c_{g,k}}{2(\omega_{h,k}^2 - \omega_{g,r}^2)} \left[(\ln p_j^2 - \ln \omega_{h,r}^2) (p_j^2 - \omega_{h,r}^2)^{-1} - (\ln p_j^2 - \ln \omega_{g,k}^2) (p_j^2 - \omega_{g,k}^2)^{-1} \right] \right\}. \quad (30b)$$

从式(30)可以看出 T_j^0, T_j^1 只与结构响应特征值和地震动功率谱密度的参数有关,该表达式为显式表达式且无需积分,为封闭解。

由式(27)及式(30a)可获得结构响应的0阶谱矩;根据式(27)及式(30b)可获得结构响应的1阶谱矩。由于 T_j^0, T_j^1 为封闭解,故0阶、1阶谱矩也为封闭解。

4.2 地震动响应的2阶谱矩简明封闭解

由随机振动理论^[24],结构响应的2阶谱矩可由该响应量变化率的0阶谱矩表示:

$$\alpha_{X,2}^2 = \alpha_{\dot{X},0}, \quad (31)$$

式中 $\dot{X} = dX/dt$ 。

由随机振动理论^[24],结构响应方差等于其0阶谱矩,结构响应变化率的方差等于其2阶谱矩:

$$\sigma_X^2 = \alpha_{X,0}, \quad \sigma_{\dot{X}}^2 = \alpha_{X,2}. \quad (32)$$

由式(15)、(28)、(31)及(32)可以看出,结构响应的0阶、1阶、2阶谱矩及方差均具有封闭解,且表达式简洁明了。

5 算例

5.1 本文计算方法的验证

以单自由度结构设置混联Ⅱ型惯容减震系统为例:结构质量为 $2.5 \times 10^6 \text{ kg}$;抗侧刚度为 $5.7 \times 10^8 \text{ N}\cdot\text{m}^{-1}$,

结构阻尼为 $6.3 \times 10^4 \text{ N}\cdot\text{s}\cdot\text{m}^{-1}$. Clough-Penzien 谱参数取值如下: $\omega_g = 15.71 \text{ rad}\cdot\text{s}^{-1}$, $\xi_g = 0.72$, $\omega_f = 0.15\omega_g$, $\xi_f = 0.72$, $S_0 = 2.317 \times 10^{-3} \text{ m}^2\cdot\text{s}^{-3}$. 惯容参数 $m_{in} = 1.2 \times 10^4 \text{ kg}$, $c_d = 1.0 \times 10^4 \text{ N}\cdot\text{s}\cdot\text{m}^{-1}$, $k_s = 1.0 \times 10^7 \text{ N}\cdot\text{m}^{-1}$.

为了验证本文方法所得谱矩的正确性, 采用虚拟激励法进行验证. 基于虚拟激励法的结构响应推导方式如下.

对式(1)引入虚拟激励, 则

$$-\omega^2 m \bar{x}(i\omega) + c i \omega \bar{x}(i\omega) + k \bar{x}(i\omega) + \bar{F}_{IS} \bar{x}(i\omega) = -m \sqrt{S_{\dot{x}_g}(\omega)} e^{i\omega t}. \quad (33)$$

对式(2)进行 Fourier 变换, 可得

$$\bar{F}_{IS}(i\omega) = k_s \bar{x}_1(i\omega) = c_d i \omega \bar{x}_2(i\omega) + m_{in}(-\omega^2) \bar{x}_2(i\omega), \quad (34a)$$

$$\bar{x}_2(i\omega) = \bar{x}(i\omega) - \bar{x}_1(i\omega). \quad (34b)$$

由式(33)可得

$$\bar{F}_{IS}(i\omega) = \frac{[c_d + m_{in}(-\omega)] \omega \bar{x}(i\omega)}{1 + i\omega c_d \frac{1}{k_s} - m_{in} \frac{\omega^2}{k_s}}. \quad (35)$$

将式(35)代入式(33)可得

$$\bar{x}(i\omega) = \frac{-m e^{i\omega t}}{-m\omega^2 + i\omega c_d + k + \rho(i\omega)} \sqrt{S_{\dot{x}_g}(\omega)}, \quad (36)$$

式中 $\rho(i\omega) = (i\omega c_d - m_{in} \omega^2) / \left(1 + i\omega c_d \frac{1}{k_s} - m_{in} \omega^2 \frac{1}{k_s}\right)$.

虚拟激励法的质点位移响应功率谱和速度响应功率谱分别为

$$S_x(\omega) = \bar{x}(i\omega) \bar{x}(i\omega)^* S_{\dot{x}_g}(\omega), \quad (37)$$

$$S_{\dot{x}}(\omega) = \omega^2 \bar{x}(i\omega) \bar{x}(i\omega)^* S_{\dot{x}_g}(\omega), \quad (38)$$

式中 $\bar{x}(i\omega)^* = \text{conj}[\bar{x}(i\omega)]$.

图3为两种方法的位移响应功率谱的对比图. 从图中可以看出, 本文得到的结果与虚拟激励法得到的结果一致, 验证了本文所提频率响应特征值函数二次正交化法和 Clough-Penzien 谱等效形式的正确性.

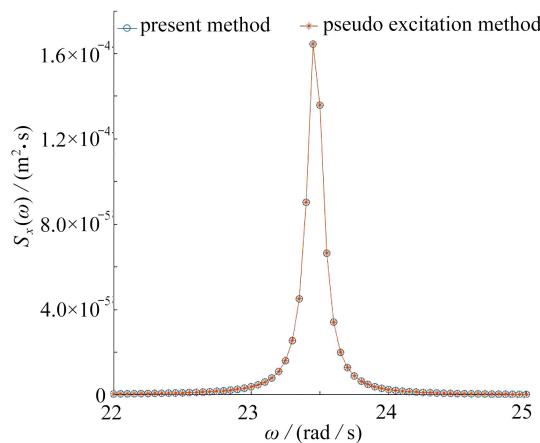


图3 结构位移功率谱
Fig. 3 The structural displacement power spectrum

为了验证谱矩的计算精度, 将本文方法与虚拟激励法(PEM)进行0~2阶谱矩的对比.

用虚拟激励法计算位移的0~2阶谱矩为

$$\lambda_0 = \int_{-\infty}^{\infty} S_x(\omega) d\omega, \quad (39a)$$

$$\lambda_1 = \int_{-\infty}^{\infty} \omega S_x(\omega) d\omega, \quad (39b)$$

$$\lambda_2 = \int_{-\infty}^{\infty} \omega^2 S_x(\omega) d\omega; \quad (39c)$$

惯容位移 x_2 的0~2阶谱矩为

$$\lambda_0 = \int_{-\infty}^{\infty} S_{x_2}(\omega) d\omega, \quad (40a)$$

$$\lambda_1 = \int_{-\infty}^{\infty} \omega S_{x_2}(\omega) d\omega, \quad (40b)$$

$$\lambda_2 = \int_{-\infty}^{\infty} \omega^2 S_{x_2}(\omega) d\omega. \quad (40c)$$

对式(39)与式(40)进行积分运算,计算结果如表1所示,当积分区间为[0, 500]时,虚拟激励法取不同的积分步矩0.01 rad/s, 0.1 rad/s, 0.5 rad/s。

表1 谱矩对比

Table 1 Comparison of spectral moments between the presented method and the virtual excitation method

calculation method	integration step Δ /(rad/s)	spectral moment		
		$\alpha_{x,0}/(10^{-5} \text{ m}^2)$	$\alpha_{x,1}/(10^{-4} \text{ m}^2/\text{s})$	$\alpha_{x,2}/(\text{m}^2/\text{s}^2)$
present method	—	7.369214319	1.721475154	0.040331515
	0.01	7.369214317	1.721475145	0.040331509
	0.1	7.247097750	1.692932655	0.039664397
PEM	0.5	1.465220153	3.434894652	0.080641421

从表1可知,积分步长越小,虚拟激励法得到的谱矩越接近本文方法的谱矩,说明了本文方法的计算精度比虚拟激励法更高。本文方法在计算0~2阶谱距所花总时间为0.009 s,虚拟激励法所花总时间分别为0.074 s,0.015 s,0.012 s,可以明显看出本文方法的计算效率更高。

为了进一步验证本文方法,将虚拟激励法与本文方法求出的惯容出力进行对比,计算结果如表2所示。当积分区间为[0, 500]时,虚拟激励法取不同的积分步矩0.5 rad/s, 0.05 rad/s, 0.01 rad/s。由表2可知,在虚拟激励法中,积分区间一定时,虚拟激励法的积分步长取的越小,越接近本文方法的结果,从而说明本文方法的正确性。

表2 惯容出力对比

Table 2 Comparison of inerter forces between the virtual excitation method and this method

calculation method	integration interval	integration step	inerter force	error
present method	[0, 500]	—	$2.466 \times 10^2 \text{ kN}$	—
		0.01 rad/s	$2.466 \times 10^2 \text{ kN}$	0
PEM	[0, 500]	0.05 rad/s	$2.466 \times 10^2 \text{ kN}$	0
		0.5 rad/s	$2.314 \times 10^2 \text{ kN}$	4.6%

为了体现惯容的减震能力,将加惯容后的结构位移与不加惯容的结构位移相比较并计算了减震率。在不加惯容的情况下,结构绝对位移为0.013435035 m,加上惯容后结构绝对位移为0.008584412 m。由此得出减震率为36.1%,说明惯容能很好地提高结构的抗震能力。

5.2 惯容参数对抗震性能的影响分析

以单自由度结构设置混联Ⅱ型惯容减震系统为例:结构质量为 $2 \times 10^4 \text{ kg}$;抗侧刚度为 $2.5 \times 10^7 \text{ N}\cdot\text{m}^{-1}$,结构阻尼为 $2.65 \times 10^4 \text{ N}\cdot\text{s}\cdot\text{m}^{-1}$ 。为研究惯容三参数对惯容减震性能的影响,利用求谱矩的方法分析惯容主要参数质量比 μ_m 、频率比 μ_ω 、阻尼比 μ_ξ 对减震性能的影响。取值范围分别为0.005~0.05, 0.5~5, 0.2~2。Clough-Penzien谱参数取值如下: $\omega_g=15.71 \text{ rad/s}$, $\xi_g=0.72$, $\omega_f=0.15\omega_g$, $\xi_f=0.72$, $S_0=2.317 \times 10^{-3} \text{ m}^2/\text{s}^3$ 。图4—6给出了上述参数的位移方差(δ)对比图,其中各点的间距为 $u_m=0.0025$, $u_\omega=0.25$, $u_\xi=0.1$ 。

由图4(a)可知,当 $\mu_\omega=1$,阻尼比不变时,结构位移方差随着质量比的增大而减小,随后将趋于平缓。同时可以看出,阻尼比越小,结构位移方差变化越平缓;在 $\mu_\xi=0.2$ 时,阻尼比太小获得结果变化不明显。由图4(b)可知,当 $\mu_\xi=1.2$,在频率不变时,结构位移方差随着质量比的增大而减小,当 $\mu_\omega=0.5$ 时随质量比变化趋势不明显。当 $\mu_\omega \geq 1.75$ 时,减震效果十分明显;当质量比一定时,频率比越大结构位移方差越小。

由图5(a)可知,当 $\mu_m=0.03$,阻尼比不变时,结构位移方差随频率先减小后增大再减小;阻尼比越大,结构

位移方差越小,在 $\mu_\omega=1$ 且 $\mu_\xi=2$ 时取得最优值。由图5(b)可知,当 $\mu_\xi=1.2$,在质量比一定时,结构位移方差随着频率比先减小后增大,在 $\mu_\omega=1$ 时取得最优值;质量比越大,结构位移方差越小。

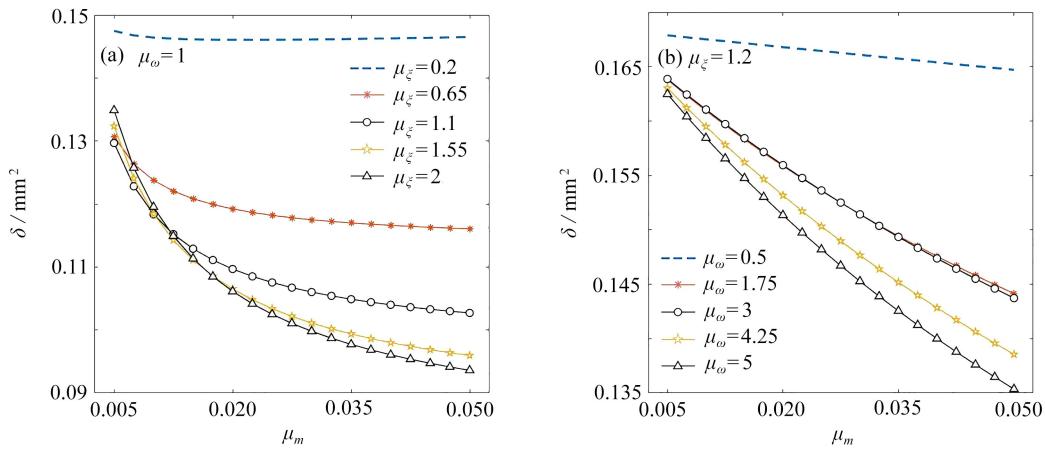


图4 惯容质量 m_{in} 对结构位移的影响
Fig. 4 Effects of inerter coefficient m_{in} on structural displacements

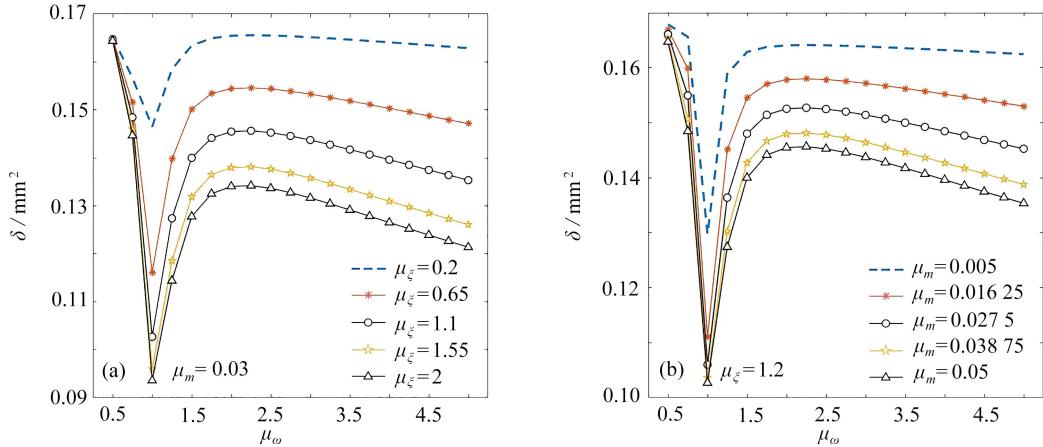


图5 惯容刚度 k_s 对结构位移的影响
Fig. 5 Effects of inerter stiffness k_s on structural displacements

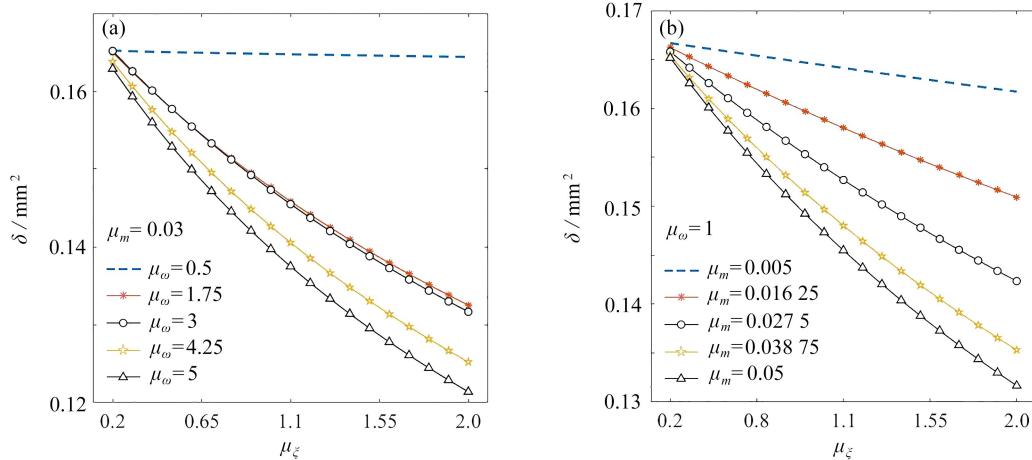


图6 惯容阻尼 c_d 对位移的影响
Fig. 6 Effects of inerter damping c_d on displacements

由图6(a)可知,当 $\mu_m=0.03$,在频率比不变时,结构位移方差随着阻尼比的增大而减小;当频率比不同时,结构位移方差随着频率比的增大而减小。在 $\mu_\omega=0.5$ 时,结构位移方差随质量变化不明显,误差较大。由图6(b)

可知,当 $\mu_\omega = 1$,在质量比一定时,位移方差随着阻尼比的增大而减小;当在同一阻尼比时,质量比越大,结构位移方差就越小.

6 结 论

本文对单自由度惯容系统混联Ⅱ型结构在Clough-Penzien谱下的平稳地震动响应进行了研究,所得结论如下:

1) 基于功率谱的二次正交化法给出了单自由度惯容系统混联Ⅱ型结构在Clough-Penzien谱下的平稳地震动响应矩的简明封闭解,通过与虚拟激励法对比验证得出本文方法在计算精度与计算效率上都优于虚拟激励法.由于采用的虚拟激励法,故本文所提方法适用于各类线性结构在Clough-Penzien谱下的平稳地震动响应矩的分析.

2) 本文的谱矩计算方法不受积分步长和积分区间的影响,计算精度比虚拟激励法显著提高,可为结构的设计和优化提供分析路径,更有利于工程的应用.

3) 通过对惯容3参数的研究,结合结构位移方差曲线变化可知:频率比不变时,随着质量比和阻尼比增大,减震性能增强;阻尼比不变时,增大质量比、频率比,减震性能增强;质量比不变时,增大频率比、阻尼比,减震性能增强;当质量比或阻尼比不变时,均在频率比为1时取得最优减震效果.

附 录

T_j^0 和 T_j^1 的推导过程如下.

$$\begin{aligned} T_j^0 &= \int_0^\infty \omega^0 A_{rj}(\omega) d\omega = \\ &S_0 \sum_{r=1}^2 \sum_{t=1}^2 c_{g,r} \int_0^\infty \frac{1}{\omega^2 + \omega_{g,r}^2} \frac{1}{p_j^2 + \omega^2} \left(\frac{1}{2} - \frac{c_{h,t}}{\omega^2 + \omega_{h,t}^2} \right) d\omega = \\ &S_0 \sum_{r=1}^2 \int_0^\infty \frac{1}{2} \frac{c_{g,r}}{\omega^2 + \omega_{g,r}^2} \frac{1}{p_j^2 + \omega^2} d\omega - \\ &S_0 \sum_{r=1}^2 \sum_{t=1}^2 \int_0^\infty \frac{c_{g,r}}{\omega^2 + \omega_{g,r}^2} \frac{c_{h,t}}{p_j^2 + \omega^2} \frac{1}{\omega^2 + \omega_{h,t}^2} d\omega. \end{aligned} \quad (\text{A1})$$

对式(A1)进行积分,可得

$$\begin{aligned} T_j^0 &= \frac{\pi}{2} \left\{ \sum_{r=1}^2 c_{g,r} [(p_j + \omega_{g,r}) \omega_{g,r} p_j]^{-1} - \right. \\ &\left. \sum_{r=1}^2 \sum_{k=1}^2 \frac{c_{h,k} c_{g,r}}{\omega_{h,k}^2 - \omega_{g,r}^2} [(\omega_{g,r} p_j^2 + \omega_{g,r}^2 p_j)^{-1} - (\omega_{h,k} p_j^2 + \omega_{h,k}^2 p_j)^{-1}] \right\} S_0, \end{aligned} \quad (\text{A2})$$

$$\begin{aligned} T_j^1 &= \int_0^\infty \omega A_{rj}(\omega) d\omega = \sum_{r=1}^2 \frac{c_{g,r}}{2} \int_0^\infty \frac{1}{\omega^2 + \omega_{g,r}^2} \frac{1}{p_j^2 + \omega^2} d\omega - \\ &\sum_{r=1}^2 \sum_{t=1}^2 \frac{c_{g,r} c_{h,t}}{2} \int_0^\infty \frac{1}{\omega^2 + \omega_{g,r}^2} \frac{1}{\omega^2 + \omega_{h,t}^2} \frac{1}{p_j^2 + \omega^2} d\omega^2 = \\ &\frac{1}{2} \left\{ \sum_{r=1}^2 c_{g,r} (\ln p_j^2 - \ln \omega_{g,r}^2) (p_j^2 - \omega_{g,r}^2)^{-1} - \right. \\ &\left. \sum_{r=1}^2 \sum_{k=1}^2 \frac{c_{h,k} c_{g,r}}{2(\omega_{h,k}^2 - \omega_{g,r}^2)} [(\ln p_j^2 - \ln \omega_{h,r}^2) (p_j^2 - \omega_{h,r}^2)^{-1} - (\ln p_j^2 - \ln \omega_{g,k}^2) (p_j^2 - \omega_{g,k}^2)^{-1}] \right\}. \end{aligned} \quad (\text{A3})$$

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