

推广的KdV方程特征问题解的存在唯一性*

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摘要

推广的 KdV 方程 $u_t + \alpha uu_x + \mu u_{x^3} + \epsilon u_{x^5} = 0$ 是典型的可积方程。它先后在研究冷等离子体中磁声波的传播^[2], 传输线中孤立波^[3]和分层流体中界面孤立波^[4]时导出。本文对推广的KdV方程的特征问题, 在 Riemann 函数的基础上, 设计一恰当结构, 并由此化特征问题为一与之等价的积分微分方程。而该积分微分方程对应的映射 E 是列自身的映射^[5], 依不动点原理, 积分微分方程有唯一的正则解, 即推广的KdV方程的特征问题有唯一解, 且由积分微分方程序列所得的迭代解于 $\bar{\Omega}$ 上一致收敛。

关键词 Riemann函数 结构 积分微分方程 不动点 一致收敛

一、设计结构

在推广的KdV方程的特征问题

$$(T) \begin{cases} u_t + \alpha uu_x + \mu u_{x^3} + \epsilon u_{x^5} = 0 & ((x,t) \in \Omega) \\ u(x,0) = h(x) \\ u(a,t) = f(t) \\ u_x(a,t) = g(t) \\ u_{x^2}(a,t) = \varphi(t) \\ u_{x^3}(a,t) = \psi(t) \\ u_{x^4}(a,t) = s(t) \end{cases}$$

中, $\bar{\Omega} = \{(x,t) : a \leq x \leq b, 0 \leq t \leq T_0, a, b, T_0 \in R^+\}$, $h(x) \in C^5[a, b]$, $f(t), g(t), \varphi(t), \psi(t), s(t) \in C^1[0, T_0]$, $h(a) = f(0), h'(a) = g(0), h''(a) = \varphi(0), h'''(a) = \psi(0), h^{(4)}(a) = s(0)$ ($a, \mu, \epsilon \in R$)

对推广的KdV方程, 若 $u \in C^{5,1}([a, b] \times [0, T_0])$, 则方程可表成算子形式

$$\begin{aligned} \mathcal{L}(u) &= \frac{1}{2} u_{x^5 t} + u_{x^5} \\ &= -\frac{1}{\epsilon} \left[u_t + \alpha uu_x + \mu u_{x^3} - \frac{\epsilon}{2} u_{x^5 t} \right] \\ &= F(u, u_x, u_t, u_{x^3}, u_{x^5 t}) \end{aligned} \tag{1.1}$$

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于是特征问题(T)化为

$$(T)' \quad \begin{cases} \mathcal{L}(u) = F(u, u_x, u_{x^2}, u_{x^3}, u_{x^4}) \\ u(x, 0) = h(x) \\ u(a, t) = f(t) \\ u_x(a, t) = g(t) \\ u_{x^2}(a, t) = \varphi(t) \\ u_{x^3}(a, t) = \psi(t) \\ u_{x^4}(a, t) = s(t) \end{cases}$$

引入方程 $\mathcal{L}(u) = 0$ 的 Riemann 函数的概念. 经计算算子 $\mathcal{L}(u)$ 的共轭算子为

$$\mathcal{L}^*(V) = \frac{1}{2}V_{x^5t} - V_{x^5} \quad (1.2)$$

对 $\forall (x_0, t_0) \in \bar{\Omega}$, 问题

$$(R) \quad \begin{cases} \mathcal{L}^*(V) = 0 \\ V(x_0, t_0; x_0, t_0) = 0 \\ V(x, t_0; x_0, t_0) = (x - x_0)^4 \\ V_x(x_0, t_0; x_0, t_0) = 0 \\ V_x(x, t_0; x_0, t_0) = 4(x - x_0)^3 \\ V_{x^2}(x_0, t_0; x_0, t_0) = 0 \\ V_{x^2}(x, t_0; x_0, t_0) = 12(x - x_0)^2 \\ V_{x^3}(x_0, t_0; x_0, t_0) = 0 \\ V_{x^3}(x, t_0; x_0, t_0) = 4!(x - x_0) \\ V_{x^4}(x_0, t_0; x_0, t_0) = 4! \exp[t - t_0] \\ V_{x^4}(x, t_0; x_0, t_0) = 4! \end{cases}$$

的解 $V(x, t; x_0, t_0) = (x - x_0)^4 \exp[t - t_0]$ 称为方程 $\mathcal{L}(u) = 0$ 的 Riemann 函数. 设

$$C_0^1(\bar{\Omega}) = \{u: u, u_x, u_t, u_{x^2}, u_{x^3}, u_{x^4} \in C(\bar{\Omega})\} \quad (1.3)$$

并于 $C_0^1(\bar{\Omega})$ 上定义范数

$$\|u\|_\lambda = \frac{\varepsilon}{4L_0} \max\{|u_t| + |uu_x| + |u_{x^2}| + |u_{x^3}| + |u_{x^4}|\} \exp[-\lambda(x_0 + t_0)] \quad (1.4)$$

其中 $L_0 = \max\{1, a, \mu, \frac{\varepsilon}{2}\}$, λ 为充分大的正常数. 可知 $C_0^1(\bar{\Omega})$ 为 Banach 空间.

对任给的 $\bar{u} \in C_0^1(\bar{\Omega})$, 讨论辅助问题

$$(T)'' \quad \begin{cases} \mathcal{L}(u) = F(\bar{u}, \bar{u}_x, \bar{u}_t, \bar{u}_{x^2}, \bar{u}_{x^3}, \bar{u}_{x^4}) \\ u(x, 0) = h(x) \\ u(a, t) = f(t) \\ u_x(a, t) = g(t) \\ u_{x^2}(a, t) = \varphi(t) \\ u_{x^3}(a, t) = \psi(t) \\ u_{x^4}(a, t) = s(t) \end{cases}$$

设计结构

$$V\mathcal{L}(u) - u\mathcal{L}^*(V) = -\frac{\partial}{\partial x} \left(\frac{1}{2}Vu_{x^4t} + Vu_{x^4} - V_x u_{x^3} + 2V_{x^2}u_{x^2} + V_{x^2}u_{x^2t} - V_x u_{x^3t} \right)$$

$$\begin{aligned}
& -\left(2V_{x^3}u_{x^2}+V_{x^2}u_{x^3}+V_{x^3}u_{x^2t}-\frac{1}{2}V_{x^2}u_{x^2t}\right) \\
& =VF(\bar{u}, \bar{u}_x, \bar{u}_t, \bar{u}_{x^2}, \bar{u}_{x^3}) \tag{1.5}
\end{aligned}$$

$$\text{设 } D = \{(x, t) : a \leq x \leq x_0 \leq b, 0 \leq t \leq t_0 \leq T_0\} \tag{1.6}$$

将(1.5)两端在 D 上积分, 且利用定解条件, 得

$$\begin{aligned}
& [u(x_0, t_0) - f(t_0)] - \exp[-t_0][h(x_0) - h(a)] \\
& + (a - x_0)[g(t_0) - \exp[-t_0]h'(a)] \\
& + (a - x_0) \int_0^{t_0} \exp[t - t_0]g(t)dt - \frac{1}{2}(a - x_0)^2[\varphi(t_0) - \exp[-t_0]\varphi(0)] \\
& - \frac{1}{2}(a - x_0)^2 \int_0^{t_0} \exp[t - t_0]\varphi(t)dt + \frac{1}{6}(a - x_0)^3[\psi(t_0) - \exp[-t_0]\psi(0)] \\
& + \frac{1}{6}(a - x_0)^3 \int_0^{t_0} \exp[-t - t_0]\psi(t)dt - \int_0^{t_0} \exp[t - t_0]f(t)dt \\
& + \int_0^{t_0} \exp[t - t_0]u(x_0, t)dt - \frac{1}{24}(a - x_0)^4[s(t_0) - \exp[-t_0]s(0)] \\
& - \frac{1}{24}(a - x_0)^4 \int_0^{t_0} \exp[t - t_0]s(t)dt \\
& = \frac{1}{12} \int_0^{t_0} \int_a^{x_0} \exp[t - t_0](x - x_0)^4 F(\bar{u}, \bar{u}_x, \bar{u}_t, \bar{u}_{x^2}, \bar{u}_{x^3}) dx dt \tag{1.7}
\end{aligned}$$

(1.7)两端关于 t_0 求导, 得

$$\begin{aligned}
& u_t(x_0, t_0) - f'(t_0) + \exp[-t_0][h(x_0) - h(a)] \\
& + (a - x_0)[g'(t_0) + \exp[-t_0]h'(a)] \\
& + (a - x_0)g(t_0) - (a - x_0) \int_0^{t_0} \exp[t - t_0]g(t)dt \\
& - \frac{1}{2}(a - x_0)^2[\varphi'(t_0) + \exp[-t_0]\varphi(0)] \\
& - \frac{1}{2}(a - x_0)^2\varphi(t_0) + \frac{1}{2}(a - x_0)^2 \int_0^{t_0} \exp[t - t_0]\varphi(t)dt \\
& + \frac{1}{6}(a - x_0)^3[\psi'(t_0) + \exp[-t_0]\psi(0)] \\
& + \frac{1}{6}(a - x_0)^3\psi(t_0) - \frac{1}{6}(a - x_0)^3 \int_0^{t_0} \exp[t - t_0]\psi(t)dt \\
& - f(t_0) + \int_0^{t_0} \exp[t - t_0]f(t)dt + u(x_0, t_0) - \int_0^{t_0} \exp[t - t_0]u(x_0, t)dt \\
& - \frac{1}{24}(a - x_0)^4[s'(t_0) + \exp[-t_0]s(0)] \\
& - \frac{1}{24}(a - x_0)^4s(t_0) + \frac{1}{24}(a - x_0)^4 \int_0^{t_0} \exp[t - t_0]s(t)dt \\
& = \frac{1}{12} \int_a^{x_0} (x - x_0)^4 F(\bar{u}, \bar{u}_x, \bar{u}_t, \bar{u}_{x^2}, \bar{u}_{x^3})|_{t=t_0} dx
\end{aligned}$$

$$-\frac{1}{12} \int_0^{t_0} \int_a^{x_0} \exp[t-t_0] (x-x_0)^4 F(\bar{u}, \bar{u}_x, \bar{u}_t, \bar{u}_{x^2}, \bar{u}_{xt}) dx dt \quad (1.8)$$

(1.7) + (1.8) 且经整理, 得

$$\begin{aligned} & u_s(x_0, t_0) + 2u(x_0, t_0) \\ &= f'(t_0) + 2f(t_0) - (a-x_0)[g'(t_0) + 2g(t_0)] + \frac{1}{2}(a-x_0)^2[\varphi'(t_0) + 2\varphi(t_0)] \\ & \quad - \frac{1}{12}(a-x_0)^3[\psi'(t_0) + 2\psi(t_0)] + \frac{1}{24}(a-x_0)^4[s'(t_0) + 2s(t_0)] \\ & \quad + \frac{1}{12} \int_a^{x_0} (x-x_0)^4 F(\bar{u}, \bar{u}_x, \bar{u}_t, \bar{u}_{x^2}, \bar{u}_{xt}) |_{t=t_0} dx \end{aligned} \quad (1.9)$$

解方程(1.9), 得

$$\begin{aligned} u(x_0, t_0) = & \exp[-2t_0] \left\{ h(x_0) + f(t_0) \exp[2t_0] \right. \\ & - f(0) - (a-x_0)[g(t_0) \exp[2t_0] - g(0)] \\ & + \frac{1}{2}(a-x_0)^2[\varphi(t_0) \exp[2t_0] - \varphi(0)] \\ & - \frac{1}{12}(a-x_0)^3[\psi(t_0) \exp[2t_0] - \psi(0)] \\ & \left. + \frac{1}{24}(a-x_0)^4[s(t_0) \exp[2t_0] - s(0)] \right. \\ & \left. + \frac{1}{12} \int_0^{t_0} \int_a^{x_0} \exp[2t] (x-x_0)^4 F(\bar{u}, \bar{u}_x, \bar{u}_t, \bar{u}_{x^2}, \bar{u}_{xt}) dx dt \right\} \end{aligned} \quad (1.10)$$

上述结果表明, 若辅助问题 $(T)''$ 有正则解 u , 则必满足公式 (1.10), 反之, 由公式 (1.10) 给出的函数 u , 不难验证, 它必为问题 $(T)''$ 的正则解.

于是有以下引理.

引理 特征问题 $(T)'$ 有正则解的充分必要条件是积分微分方程

$$\begin{aligned} u(x_0, t_0) = & \exp[-2t_0] \left\{ h(x_0) + f(t_0) \exp[2t_0] \right. \\ & - f(0) - (a-x_0)[g(t_0) \exp[2t_0] - g(0)] \\ & + \frac{1}{2}(a-x_0)^2[\varphi(t_0) \exp[2t_0] - \varphi(0)] \\ & - \frac{1}{6}(a-x_0)^3[\psi(t_0) \exp[2t_0] - \psi(0)] \\ & + \frac{1}{24}(a-x_0)^4[s(t_0) \exp[2t_0] - s(0)] \\ & \left. + \frac{1}{12} \int_0^{t_0} \int_a^{x_0} \exp[2t] (x-x_0)^4 F(u, u_x, u_t, u_{x^2}, u_{xt}) dx dt \right\} \end{aligned} \quad (1.11)$$

有解; 若有解, 则它们的解相同.

二、解的存在唯一性

在 $C_0^1(\bar{Q})$ 上定义映射

$$\begin{aligned}
[Eu](x_0, t_0) = & \exp[-2t_0] \left\{ h(x_0) + f(t_0) \exp[2t_0] \right. \\
& - f(0) - (a-x_0)[g(t_0) \exp[2t_0] - g(0)] \\
& + \frac{1}{2} (a-x_0)^2 [\varphi(t_0) \exp[2t_0] - \varphi(0)] \\
& - \frac{1}{6} (a-x_0)^3 [\psi(t_0) \exp[2t_0] - \psi(0)] \\
& + \frac{1}{24} (a-x_0)^4 [s(t_0) \exp[2t_0] - s(0)] \\
& \left. + \frac{1}{12} \int_0^{t_0} \int_a^{x_0} \exp[2t] (x-x_0)^4 F(u, u_x, u_t, u_{x^3}, u_{x^5 t}) dx dt \right\} \quad (2.1)
\end{aligned}$$

若证明了 E 为由 $C_0^5(\bar{\Omega})$ 到自身的映射, 则特征问题 (T') 就化成了不动点问题.

定理 特征问题 $(T)'$ 有唯一的正则解 $u^* \in C_0^5(\Omega)$, 且对给的 $u_0 \in C_0^5(\bar{\Omega})$, 序列

$$\begin{aligned}
u_n(x_0, t_0) = & \exp[-2t_0] \left\{ h(x_0) + f(t_0) \exp[2t_0] \right. \\
& - f(0) - (a-x_0)[g(t_0) \exp[2t_0] - g(0)] \\
& + \frac{1}{2} (a-x_0)^2 [\varphi(t_0) \exp[2t_0] - \varphi(0)] \\
& - \frac{1}{6} (a-x_0)^3 [\psi(t_0) \exp[2t_0] - \psi(0)] \\
& + \frac{1}{24} (a-x_0)^4 [s(t_0) \exp[2t_0] - s(0)] \\
& \left. + \frac{1}{12} \int_0^{t_0} \int_a^{x_0} \exp[2t] (x-x_0)^4 F(u_{n-1}, u_{n-1 x}, u_{n-1 t}, u_{n-1 x^3}, u_{n-1 x^5 t}) dx dt \right\} \quad (2.2)
\end{aligned}$$

于 $\bar{\Omega}$ 上一致收敛于 u^* .

证明 证明 E 为由 $C_0^5(\bar{\Omega})$ 到自身的映射.

对 $\forall u, \bar{u} \in C_0^5(\bar{\Omega})$, 设

$$\begin{aligned}
M = & \max_{(x_0, t_0) \in \bar{\Omega}} \left\{ 4 \left| \int_0^{t_0} \int_a^{x_0} (x-x_0)^4 \exp[2(t-t_0)] F(u, u_x, u_t, u_{x^3}, u_{x^5 t}) dx dt \right| \right\} \\
N = & \max_{(x_0, t_0) \in \bar{\Omega}} \left\{ 4 \left| \int_0^{t_0} \int_a^{x_0} (x-x_0)^3 \exp[2(t-t_0)] F(\bar{u}, \bar{u}_x, \bar{u}_t, \bar{u}_{x^3}, \bar{u}_{x^5 t}) dx dt \right| \right\} \\
R(x_0, t_0) = & \exp[-2t_0] \{ h(x_0) + f(t_0) \exp[2t_0] \\
& - f(0) - (a-x_0)[g(t_0) \exp[2t_0] - g(0)] \\
& + \frac{1}{2} (a-x_0)^2 [\varphi(t_0) \exp[2t_0] - \varphi(0)] \\
& - \frac{1}{6} (a-x_0)^3 [\psi(t_0) \exp[2t_0] - \psi(0)] \\
& + \frac{1}{24} (a-x_0)^4 [s(t_0) \exp[2t_0] - s(0)] \\
G(x_0, t_0) = & \exp[-2t_0] \{ h'(x_0) + [g(t_0) \exp[2t_0] - g(0)] \\
& - (a-x_0)[\varphi(t_0) \exp[2t_0] - \varphi(0)]
\end{aligned}$$

$$+ \frac{1}{2} (a-x_0)^2 [\psi(t_0) \exp[2t_0] - \psi(0)]$$

$$- \frac{1}{6} (a-x_0)^3 [s(t_0) \exp[2t_0] - s(0)]$$

$$\bar{R} = \max_{(x_0, t_0) \in \bar{\Omega}} |R(x_0, t_0)|, \quad \bar{G} = \max_{(x_0, t_0) \in \bar{\Omega}} |G(x_0, t_0)|$$

于是

$$|[Eu]_t - [E\bar{u}]_t| \leq \frac{(b-a)^4}{4\lambda} \left(\frac{1}{12} + \frac{1}{6\lambda} \right) \|u - \bar{u}\|_\lambda \exp[\lambda(x_0 + t_0)]$$

$$|[Eu][Eu]_x - [E\bar{u}][E\bar{u}]_x| \leq \frac{1}{48\lambda^2} [(a-b)^4 \bar{G} + (b-a)^3 \bar{R} + (b-a)^3 M + (b-a)^4 N] \cdot \|u - \bar{u}\|_\lambda \exp[\lambda(x_0 + t_0)]$$

$$|[Eu]_{x^3} - [E\bar{u}]_{x^3}| \leq \frac{1}{2\lambda^2} (b-a) \|u - \bar{u}\|_\lambda \exp[\lambda(x_0 + t_0)]$$

$$|[Eu]_{x^5 t} - [E\bar{u}]_{x^5 t}| \leq \frac{1}{2\lambda} \|u - \bar{u}\|_\lambda \exp[\lambda(x_0 + t_0)]$$

$$+ \frac{1}{2} \|u - \bar{u}\|_\lambda \exp[\lambda(x_0 + t_0)]$$

依(1.4), 则 $\|Eu - E\bar{u}\|_\lambda \leq \left(\frac{1}{2} + \frac{I}{\lambda} \right) \|u - \bar{u}\|_\lambda$

其中
$$I = \frac{1}{4} \left\{ (b-a)^4 \left(\frac{1}{12} + \frac{1}{6\lambda} \right) + \frac{1}{12\lambda} [(a-b)^4 \bar{G} + (b-a)^3 \bar{R} + (b-a)^3 M + (b-a)^4 N] + \frac{2}{\lambda} (b-a) + 2 \right\}$$

选择 λ , 使 $\frac{I}{\lambda} < \frac{1}{2}$, 则 E 为由 $C_0^5(\bar{\Omega})$ 到自身的映射. 依不动点原理. 可知映射 E 有唯一的不动点 $u^* \in C_0^5(\bar{\Omega})$, 即积分微分方程(1.11)有唯一正则解 u^* . 依引理. 特征问题 $(T)'$ 有唯一正则解 u^* .

$$\text{令 } M_{n+p-1} = \max_{(x_0, t_0) \in \bar{\Omega}} \left\{ 4 \left| \int_0^{t_0} \int_a^{x_0} \exp[2(t-t_0)] (x-x_0)^4 \cdot F(u_{n+p-1}, u_{n+p-1}, x, u_{n+p-1}, t, u_{n+p-1}, x^3, u_{n+p-1}, x^5) dx dt \right| \right\}$$

$$N_{n-1} = \max_{(x_0, t_0) \in \bar{\Omega}} \left\{ 4 \left| \int_0^{t_0} \int_a^{x_0} \exp[2(t-t_0)] (x-x_0)^3 \cdot F(u_{n-1}, u_{n-1}, x, u_{n-1}, t, u_{n-1}, x^3, u_{n-1}, x^5) dx dt \right| \right\}$$

同于上述证明, 有

$$\begin{aligned} \|u_{n+p} - u_n\|_\lambda &\leq \left(\frac{1}{2} + \frac{I_{n-1}}{\lambda} \right) \|u_{n+p-1} - u_{n-1}\|_\lambda \\ &\leq \left(\frac{1}{2} + \frac{I_{n-1}}{\lambda} \right) \left(\frac{1}{2} + \frac{I_{n-2}}{\lambda} \right) \|u_{n+p-2} - u_{n-2}\|_\lambda \\ &\leq \dots \end{aligned}$$

$$\leq \prod_{k=1}^n \left(\frac{1}{2} + \frac{I_{n-k}}{\lambda} \right) \|u_p - u_0\|_\lambda$$

选择 λ , 使 $\frac{I_{n-k}}{\lambda} < \frac{1}{2}$ ($k=1, 2, 3, \dots, n$), 于是

$$\prod_{k=1}^n \left(\frac{1}{2} + \frac{I_{n-k}}{\lambda} \right) \xrightarrow{n \rightarrow \infty} 0$$

即 $\|u_{n+p} - u_n\|_\lambda \xrightarrow{n \rightarrow \infty} 0$, 依 Cauchy 准则, 则由积分微分方程(2.2)所确定的特征问题(T)' 的迭代解于 $\bar{\Omega}$ 上一致收敛于 u^* .

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The Uniqueness and Existence of Solution of the Characteristic Problem on the Generalized KdV Equation

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Abstract

The generalized KdV equation, $u_t + auu_x + \mu u_x^3 + eu_{x^5} = 0$, is a typical integrable equation. It is derived by studying the disseminations of magnet sound wave in cold plasma[2], the isolated wave in transmission line[3], and the isolated wave in the boundary surface of the divided layer fluid [4]. For the characteristic problem of the generalized KdV equation, this paper, on the basis of Riemann function, designs a suitable structure, then changes the characteristic problem into an equivalent integral and differential equation whose corresponding mapping E is a mapping to itself[5]. According to the principle of fixed point, the above integral and differential equation has a unique regular solution, so the characteristic problem of the generalized KdV equation has a unique solution. The iterative solution derived from the integral-differential equation sequence is uniformly convergent in $\bar{\Omega}$.

Key words Riemann function, structure, integral and differential equation, fixed point, uniformly convergent